

Ambarzumian Theorem for Quantum Graphs with Magnetic Potential

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We consider a compact **metric graph** $\Gamma = (\mathbf{E}, \mathbf{V})$

- Edges e_1, \dots, e_N , $e_n = [x_{2n-1}, x_{2n}]$ and $\ell_n = x_{2n} - x_{2n-1}$.
- Vertices v_1, \dots, v_M are a partition of the endpoints x_1, \dots, x_{2N}
- The total length is $\ell(\Gamma) = \sum_{n=1}^N \ell_n$
- $L^2(\Gamma) = \bigoplus_{n=1}^N L^2(e_n)$ Hilbert space of functions on Γ .

Take $a \in \bigoplus C^1(e_n)$, $q \in \bigoplus L^\infty(e_n)$. The **Magnetic Schrödinger operator** $L_{q,a}^{st}$ acts as

$$-\left(\frac{d}{dx} - ia\right)^2 + q$$

Let $\partial u(x_j) = (-1)^{j+1} \lim_{x \rightarrow x_j} (u'(x) - ia(x)u(x))$. Then $\text{dom}(L_{q,a}^{st})$ consists of $W^{2,2}$ functions with **Standard conditions**:

$$\begin{cases} u(x_{j_1}) = u(x_{j_2}), & x_{j_1}, x_{j_2} \in v_m \\ \sum_{x_j \in v_m} \partial u(x_j) = 0 \end{cases}$$

We have an **Ambarzumian theorem** for Schrödinger operators on metric graphs

Theorem (Davies, 2013). *If the lowest eigenvalue $\lambda_1(L_q^{st}) \geq 0$ and $\limsup_{n \rightarrow \infty} \lambda_n(L_q^{st}) - \lambda_n(L_0^{st}) \leq 0$ then $q \equiv 0$*

The magnetic potential

Let Θ be a function s.t. $(\Theta|_{e_n})' = a|_{e_n}$.

Then $(\frac{d}{dx} - ia)e^{i\Theta} = e^{i\Theta} \frac{d}{dx}$

$L_q^\Theta := e^{-i\Theta} L_{q,a}^{st} e^{i\Theta}$ acts, on each edge, as $-\frac{d^2}{dx^2} + q$ with vertex conditions

$$\begin{cases} e^{i\Theta(x_{j_1})} u(x_{j_1}) = e^{i\Theta(x_{j_2})} u(x_{j_2}), & x_{j_1}, x_{j_2} \in v^m \\ \sum_{x_j \in v^m} e^{i\Theta(x_j)} \partial u(x_j) = 0, & x_j \in v^m \end{cases}$$

where $\partial u(x_j) = (-1)^{j+1} \lim_{x \rightarrow x_j} u'(x)$

If $L_q^\Theta = L_q^{st}$ then we say that the magnetic potential is **globally eliminated**. The following are equivalent:

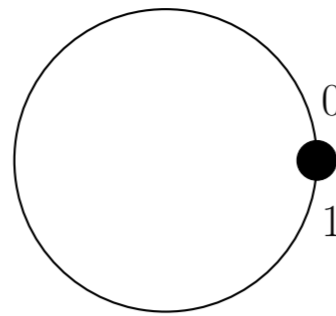
- The magnetic potential can be globally eliminated
- $e^{i\Theta}$ is continuous on Γ for some Θ
- For any cycle C : $\oint_C a dx \in 2\pi\mathbb{Z}$

For a magnetic Laplacian $L_{0,a}^{st}$ the lowest eigenvalue is given by

$$\lambda_1(L_{0,a}^{st}) = \min_{\substack{u \in C(\Gamma) \\ u \in W^{1,2}}} \frac{\int_\Gamma |u' - iau|^2 dx}{\int_\Gamma |u|^2 dx},$$

Proposition. $\lambda_1(L_{0,a}^{st}) = 0$ if and only if the magnetic potential can be globally eliminated

Proof. We have continuous u such that $\int_\Gamma |u' - iau|^2 dx = 0$. Then u must be of the form $u = e^{i\Theta}$ where $\Theta' = a$. \square



Example Let $L_{0,a}^{st}$ be an operator on a loop Γ consisting of one edge $[0, 1]$. Let $\Theta(x) = \int_0^x a(t) dt$.

- The eigenvalues of $L_{0,a}^{st}$ are $(2\pi n + \Theta(1))^2$ where $n \in \mathbb{Z}$.
- $L_{0,a}^{st}$ is isospectral to $L_{0,0}^{st}$ if and only if $\Theta(1) = 2\pi n$ for $n \in \mathbb{Z}$
- If $\Theta(1) = 2\pi n$ then $e^{i\Theta(x)} \in C(\Gamma)$ and $e^{-i\Theta} L_{0,a}^{st} e^{i\Theta} = L_{0,0}^{st}$

Reference Laplacians

The reference Laplacian of a Schrödinger operator L_q (with arbitrary vertex conditions) on Γ is the operator with quadratic form $Q(u) = \int_\Gamma |u'|^2 dx$, $\text{dom}(Q) = \overline{\text{dom}(L_q)}$.

Theorem (Kurasov, Suhr, 2020). *If L_{q_1}, L_{q_2} are two Schrödinger operators on Γ_1 and Γ_2 respectively and*

$$\sqrt{\lambda_n(L_{q_1})} - \sqrt{\lambda_n(L_{q_2})} \rightarrow 0 \text{ as } n \rightarrow \infty$$

then their reference Laplacians are isospectral.

For two magnetic Schrödinger operators $L_{q_1, a_1}^{st}, L_{q_2, a_2}^{st}$ if

$$\sqrt{\lambda_n(L_{q_1, a_1}^{st})} - \sqrt{\lambda_n(L_{q_2, a_2}^{st})} \rightarrow 0 \text{ as } n \rightarrow \infty$$

then L_{0, a_1}^{st} is isospectral to L_{0, a_2}^{st} .

Magnetic Ambarzumian Theorem

Theorem. *If $\lambda_1(L_{q,a}^{st}) \geq 0$ and $\lambda_n(L_{q,a}^{st}) - \lambda_n(L_{0,0}^{st}) \rightarrow 0$ then $q \equiv 0$ and a can be globally eliminated.*

Proof. If $\lambda_n(L_{q,a}^{st}) - \lambda_n(L_{0,0}^{st}) \rightarrow 0$ then $L_{0,a}^{st}$ and $L_{0,0}^{st}$ are isospectral.

Then $\lambda_1(L_{0,a}^{st}) = \lambda_1(L_{0,0}^{st}) = 0$ so the magnetic potential can be globally eliminated.

Since a can be eliminated $L_{q,0}^{st}$ is isospectral to $L_{q,a}^{st}$.

Then by the theorem of Davies $q \equiv 0$. \square

Generalizations

There exists several other Ambarzumian type theorems such as

1. Nicaise (1987): For an operator L_0^{st} on Γ , $\lambda_2 = (\frac{\pi}{\ell(\Gamma)})^2$ iff Γ is an interval
2. Boman, Kurasov, Suhr (2018): If $L_{q,a}^{st}$ on Γ is isospectral to L_0^{st} on an interval then $q \equiv 0$ and Γ is an interval
3. Bifulco and Kerner (2024): For certain operators* L_q^σ with delta coupling coefficients σ : If L_q^σ is isospectral to L_0^{st} then $q \equiv 0$ and $\sigma = 0$

These results can also be extended to include magnetic potential. For (1) we consider an operator $L_{0,a}^{st}$. If $\lambda_1 = 0$, $\lambda_2 = (\frac{\pi}{\ell(\Gamma)})^2$ then a can be eliminated and Γ is an interval. For (2) we assume $L_{q,a}^{st}$ is isospectral to a Laplacian on an interval. Then a can be eliminated $q \equiv 0$ and Γ is an interval. For (3) we assume $L_{q,a}^\sigma$ is isospectral to $L_{0,0}^{st}$. Then a can be eliminated $q \equiv 0$ and $\sigma = 0$.

*We require the operator to satisfy any of the 3 following conditions:

$$\sum_m \sigma_m \left(1 - \frac{2}{\deg(v^m)}\right) \leq 0, \text{ or } \int_\Gamma q dx \geq 0 \text{ and all } \sigma_m \geq 0, \\ \square \text{ or } \left(\frac{1}{2} \min_m \deg(v^m) - 1\right) \int_\Gamma q dx \geq 0 \text{ and all } \sigma_m \leq 0.$$